

# WIDG Journal Club

## Cosmic Axion Spin Precession Experiment

B. Brubaker    A. Malagon

Yale University

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# Overview

## Cosmic Axion Spin Precession Experiment (CASPER)

Dmitry Budker,<sup>1,2</sup> Peter W. Graham,<sup>3</sup> Micah Ledbetter,<sup>4</sup> Surjeet Rajendran,<sup>3</sup> and Alex Sushkov<sup>5</sup>

<sup>1</sup>*Department of Physics, University of California, Berkeley, California 94720, USA*

<sup>2</sup>*Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

<sup>3</sup>*Stanford Institute for Theoretical Physics, Department of Physics, Stanford University, Stanford, CA 94305*

<sup>4</sup>*AOSense, 767 N. Mary Ave, Sunnyvale, CA, 94085-2909*

<sup>5</sup>*Department of Physics, and Department of Chemistry and Chemical Biology,  
Harvard University, Cambridge, MA 02138, USA.*

We propose an experiment to search for QCD axion and axion-like-particle (ALP) dark matter. Nuclei that are interacting with the background axion dark matter acquire time-varying CP-odd nuclear moments such as an electric dipole moment. In analogy with nuclear magnetic resonance, these moments cause precession of nuclear spins in a material sample in the presence of a background electric field. This precession can be detected through high-precision magnetometry. With current techniques, this experiment has sensitivity to axion masses  $m_a \lesssim 10^{-9}$  eV, corresponding to theoretically well-motivated axion decay constants  $f_a \gtrsim 10^{16}$  GeV. With improved magnetometry, this experiment could ultimately cover the entire range of masses  $m_a \lesssim \mu\text{eV}$ , just beyond the region accessible to current axion searches.

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- Overview of parameter space and current detection strategies: often unfavorable scaling with small coupling constants
- One measurable effect: dark matter axion field induces oscillating EDMs in nucleons
- CASPER plans to search for this effect using NMR techniques.



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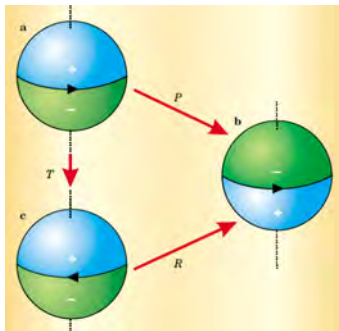
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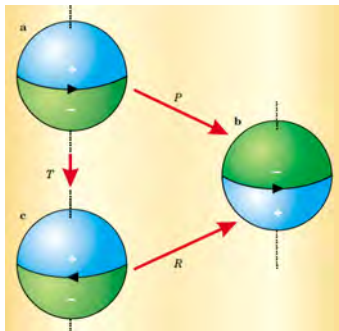
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- QCD Lagrangian contains *CP*-violating term  $\mathcal{L}_\theta \propto \theta G\tilde{G}$
- $\theta$  = sum of free parameters from strong and electroweak sectors: theory suggests  $\theta \sim \mathcal{O}(1)$ , but experiments constrain  $\theta < 10^{-10}$   
 $\Rightarrow$  *The Strong CP Problem*

# Symmetry Violation and EDMs

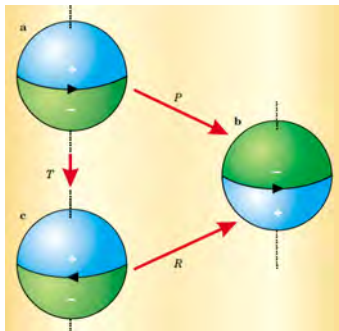


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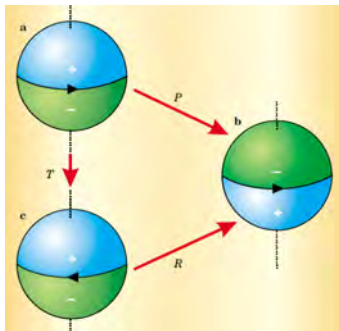
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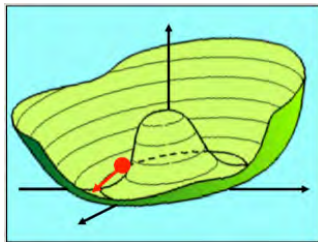
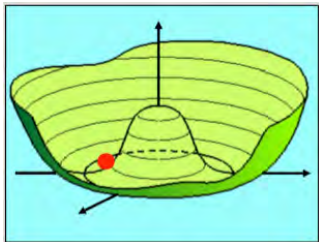
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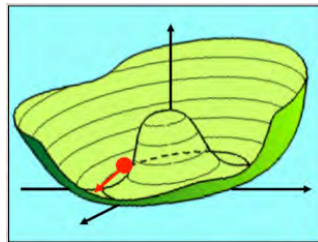
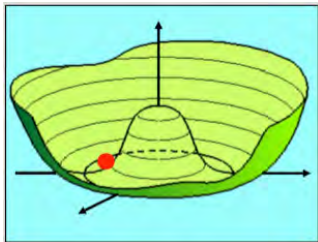
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- Magnetic and Electric Dipole Moments transform oppositely under  $P$  and  $T$  reversal
- $CP$  violation in QCD  $\Rightarrow$  neutron EDM. Non-observation constrains  $\theta < 10^{-10}$  rad



# Solving the Strong $CP$ Problem: Axions

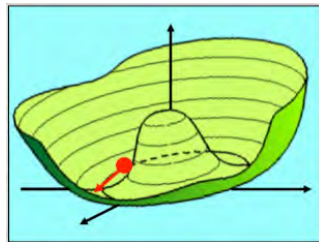
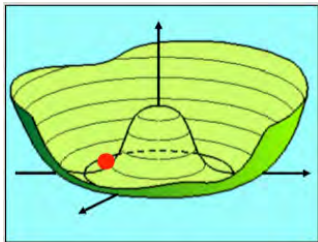


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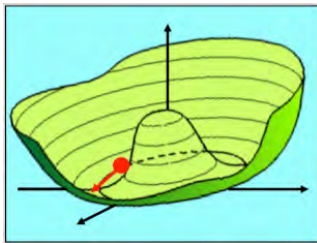
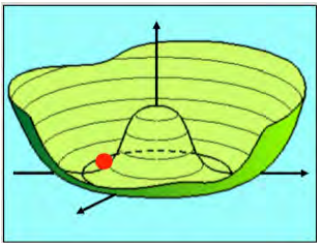
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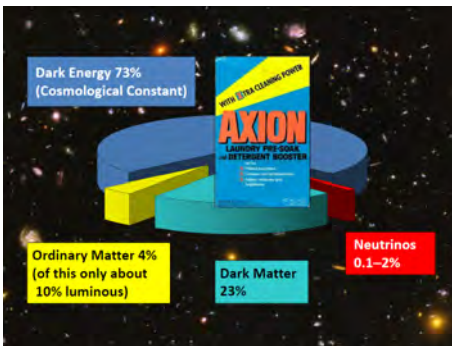


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- QCD instanton effects tilt potential:  $\theta \rightarrow 0$  dynamically. Zero-point fluctuations w/ mass  $m_a \sim \Lambda_{\text{QCD}}^2/f_a \Leftrightarrow$  axions

# Axion Dark Matter

$$\mathcal{L} = \frac{1}{2} (\partial^\mu a)^2 - \frac{1}{2} m_a^2 a^2 + C \frac{a}{f_a} G\tilde{G} + C' \frac{a}{f_a} F\tilde{F} + \dots$$

- Couplings scale as  $g \propto f_a^{-1} \propto m_a$
- $\Rightarrow$  Low-mass axions ( $m_a < 1$  meV) could be dark matter!



# Properties of Axion Dark Matter

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- Axions have virial velocity  $v \sim 10^{-3}$ 
  - Coherence length  $\lambda_a \sim \frac{1}{m_a v} > 1 \text{ m} \Rightarrow$  spatially uniform on laboratory scales
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- Solving E-L equations without couplings:  $a = a_0 \cos(m_a t)$  with  $\rho_{\text{DM}} = \frac{1}{2} m_a^2 a_0^2$

# Axion-Photon Coupling

Most axion detection schemes use axion-photon coupling

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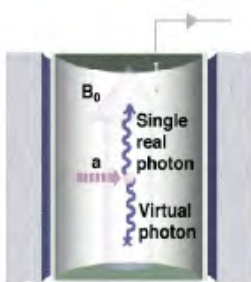
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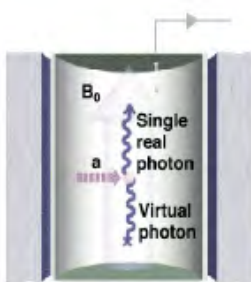
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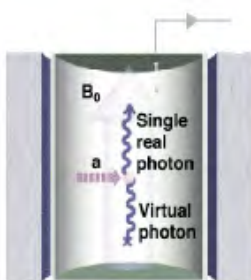
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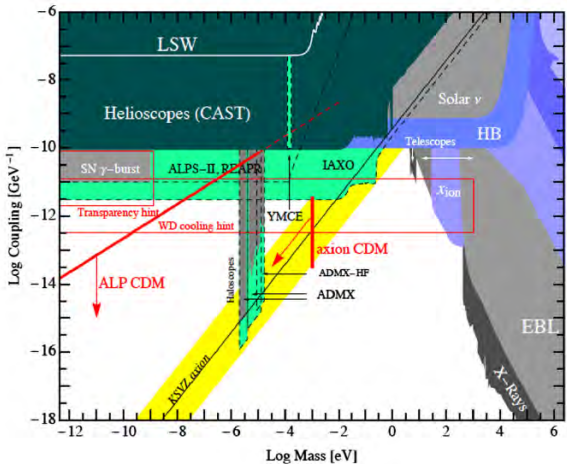
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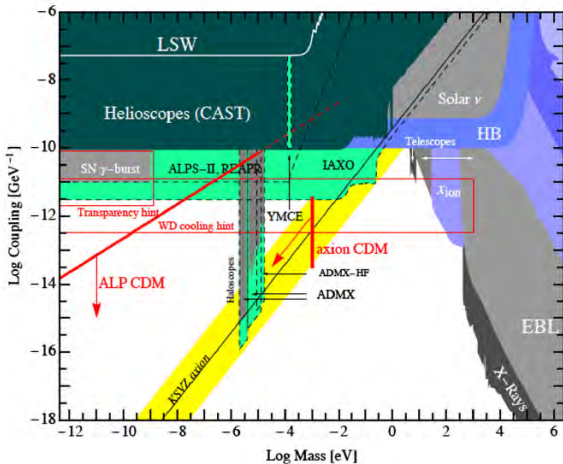
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- Resonantly enhanced in cavity of size  $L \approx \lambda_\gamma = m_a^{-1}$
- Conversion power  $P \sim g_{a\gamma\gamma}^2 \Leftarrow$  unfavorable scaling, even worse without an axion source!

# Axion Parameter Space



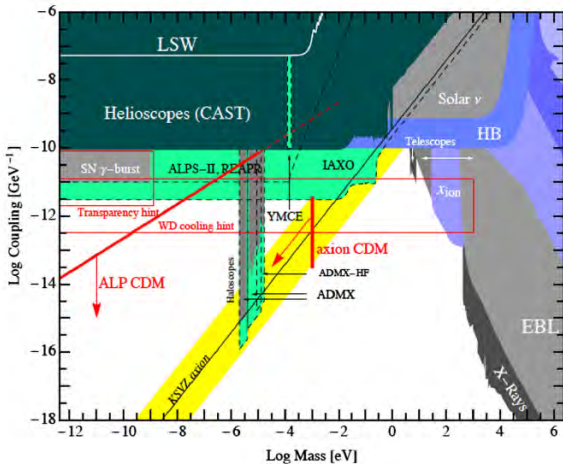
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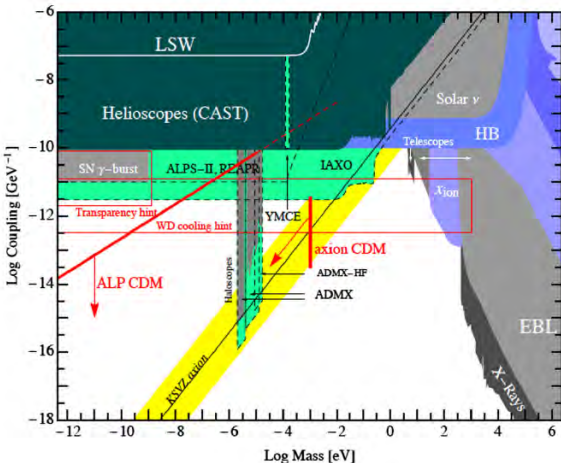
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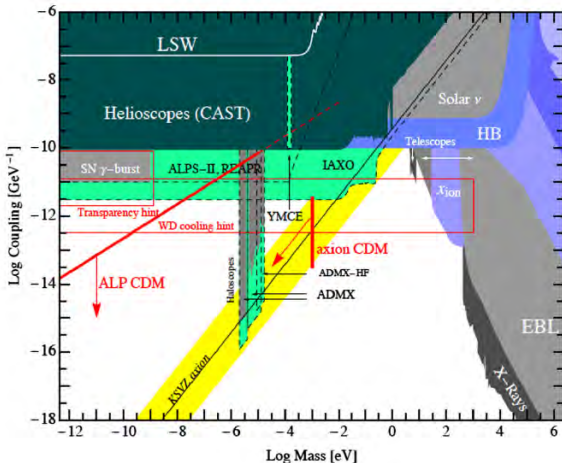


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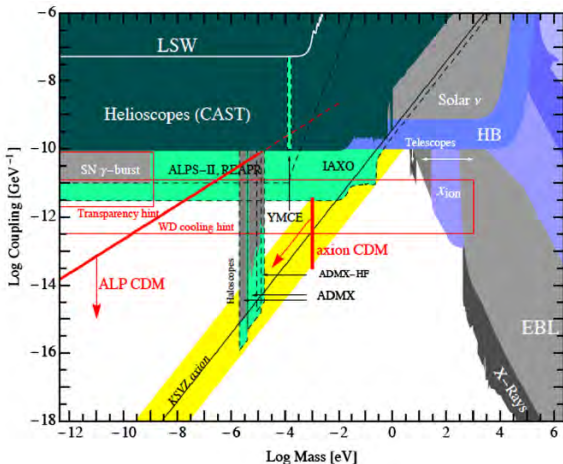
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- ALPs: No relation between mass and coupling

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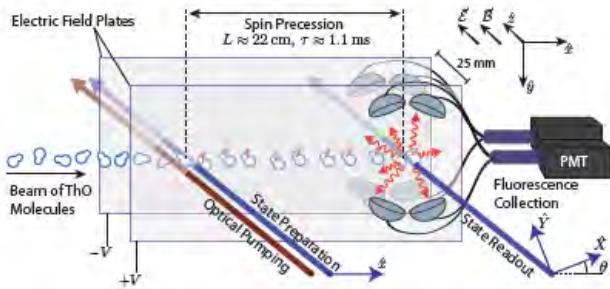
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- Observable effect linear in  $g_d$ !

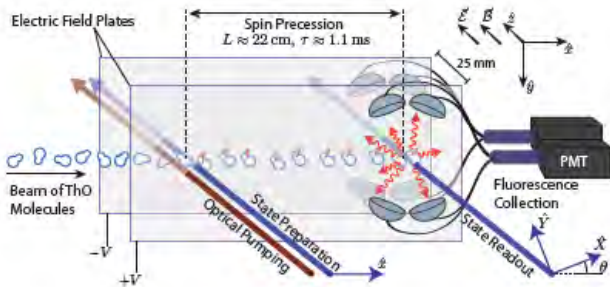
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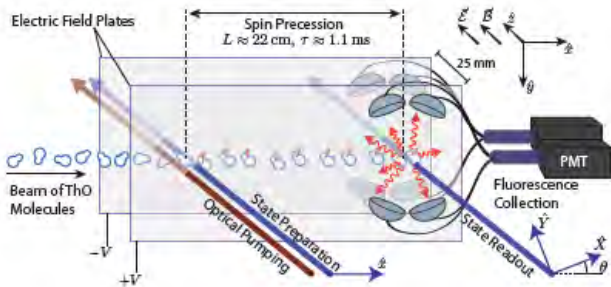


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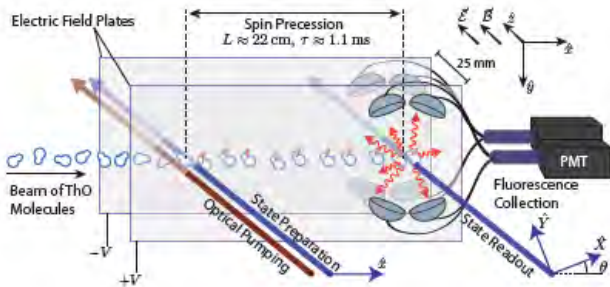
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- $t \sim 1$  ms  $\Rightarrow$  AC EDM signal averages out.

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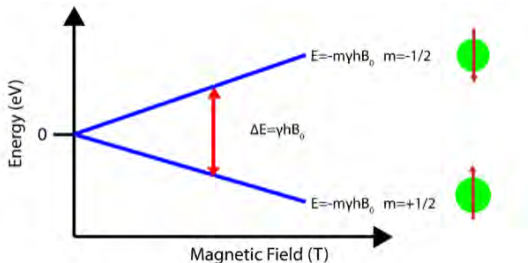
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  - Only need to precisely control high-frequency Fourier components of applied field
  - Resonant Detection is possible  $\Rightarrow$  NMR

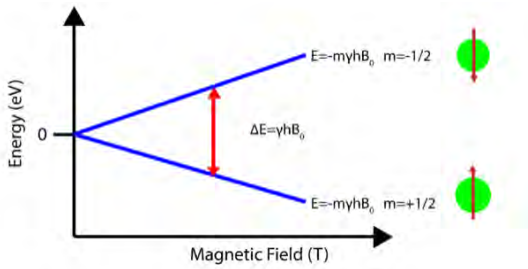
# NMR Basics – Quantum Picture



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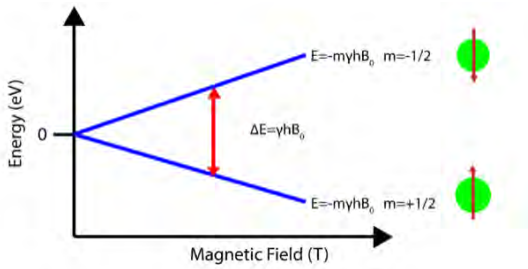


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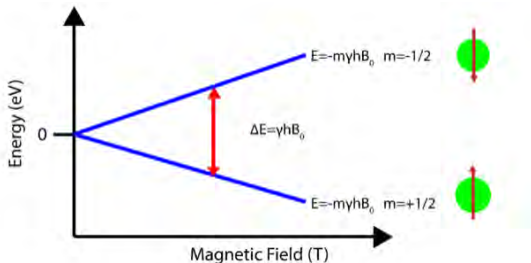
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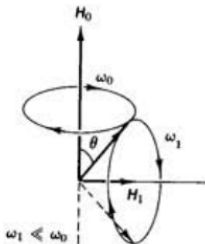
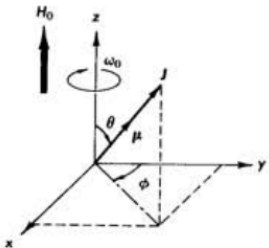
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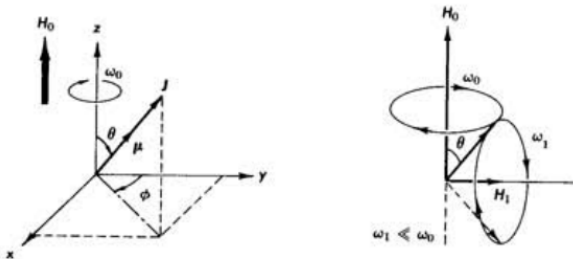


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- AC Transverse field at Larmor Frequency:  $H_1 = \gamma B_{\perp} \cos(\omega_L t)$  stimulates transitions

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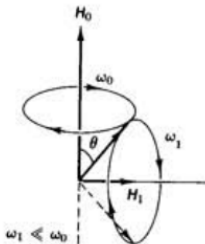
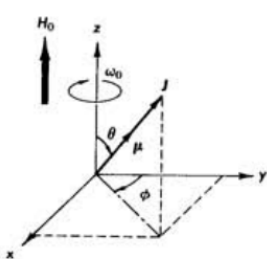


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- When  $B_{\perp}$  applied at Larmor frequency, they precess around  $B_{\perp}$  instead.

# Pulsed NMR

- In most experiments,  $B_{\perp}$  is strong, and pulsed to rotate spins into x-y plane  $\Rightarrow$  “ $\pi/2$  pulse”
- Use coil to detect transverse magnetization over time  $T_2$ .

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  - Hours at  $T \sim 4$  K: shorter is better *for pulsed NMR*



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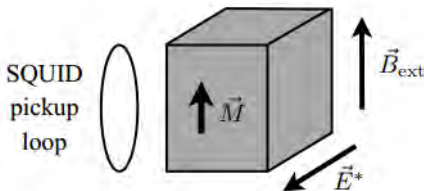
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- Perturbative  $H_1 \Rightarrow$  no saturation even for long  $T_1$ , so longer  $T_1$  is better: allows use of higher-than-thermal polarizations!

# CASPER: Experimental Setup

- Apply  $B_0$  and  $E_{\perp}$  fields to a non-paramagnetic insulator.
- Measure transverse magnetization with sensitive magnetometer (SQUID or SERF).
- Incrementally adjust  $B_0$  to scan axion mass range.



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- $np\gamma$ : Initial magnetization along  $\hat{z}$

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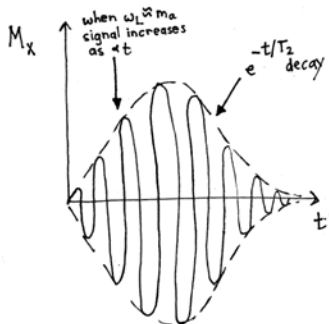
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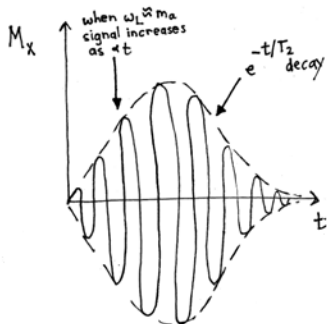
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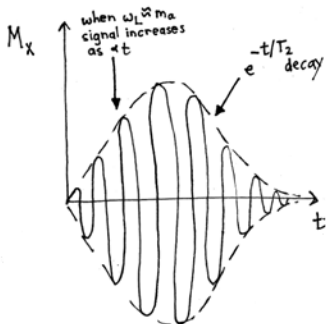
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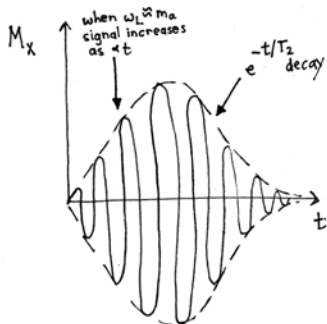
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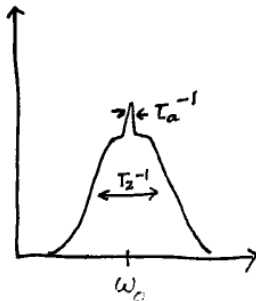
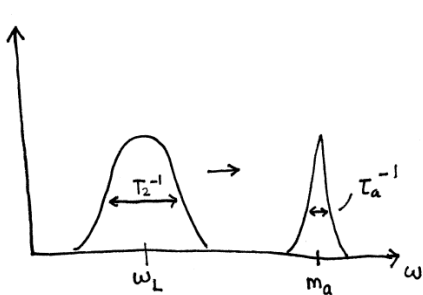
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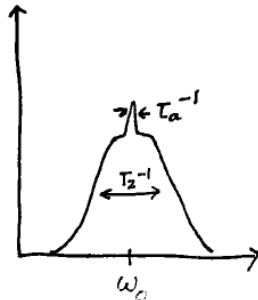
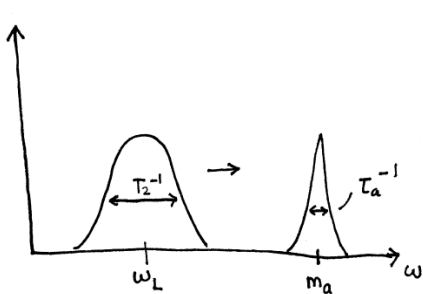
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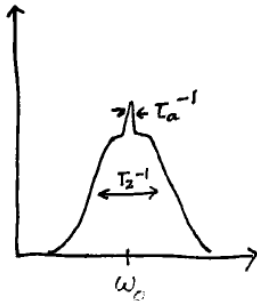
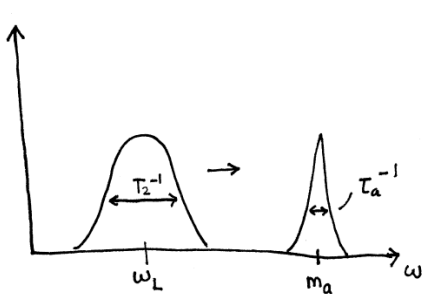
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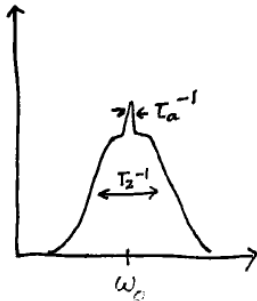
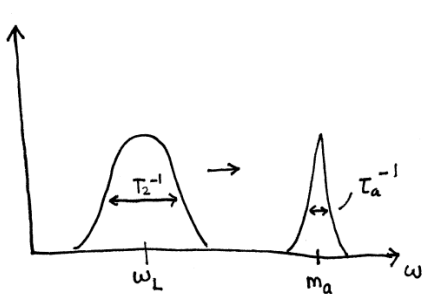
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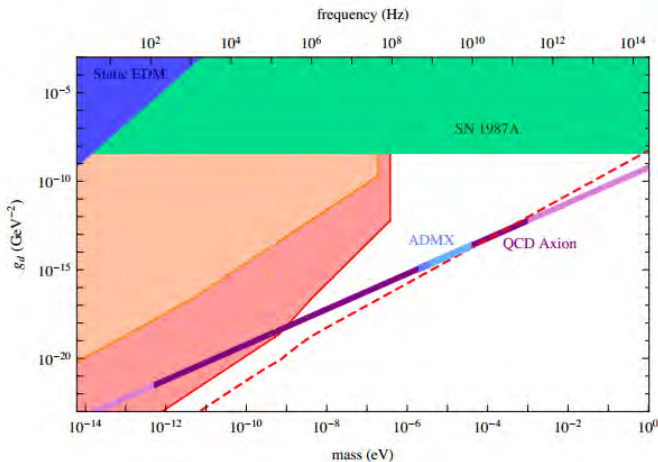


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- Accessible parameter space: anywhere below 10 T  $\Leftrightarrow$  450 MHz  $\sim 1 \mu\text{eV}$



# Proposed Limits



# Summary

- A new idea to look for axions: look for the axion-induced EDM.
- The induced EDM is small, but could be detected using NMR techniques and sensitive magnetometers.
- Observable effect scales linearly with  $g_d$
- Challenges
  - Precisely controlling transverse field fluctuations?
  - Lots of parameter space to cover
- Questions?

## Figures from

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- ACME Collaboration, Order of Magnitude Smaller Limit on the Electric Dipole Moment of the Electron, arXiv preprint, 2013.
- [http://en.wikipedia.org/wiki/Nuclear\\_magnetic\\_resonance](http://en.wikipedia.org/wiki/Nuclear_magnetic_resonance)

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